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Department of Mathematics; Computer Science and Engineering,  
Graphic Era (Deemed to be University), Dehradun, INDIA.

e-mail: [eicijmems@ijmems.in](mailto:eicijmems@ijmems.in)

Research Areas - *Operations Research, Reliability Theory, Applied Mathematics.*

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Department Mechanical Engineering,  
University of Aveiro, Campus Santiago, Aveiro, PORTUGAL.

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Centre for Reliability and Risk Management,  
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ZHAW School of Engineering,  
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Research Areas - *Fault Detection, Machine Learning, Reliability Engineering.*

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Research Areas - *Quality & Reliability Engineering, Robust Product/Process Design, Concurrent Engineering, TQM, Lean Manufacturing, Six Sigma Methodologies, Production & Operations Management, Optimization Techniques, Supply-Chain Management, Fuzzy Logic & Neural Networks, Quantitative Analysis of Operations Management.*

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Department of Mathematics,  
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Research Areas - *Survival Analysis & Semi Parametric Statistics, Medical Statistics, Model Selection, Biostatistics.*

**Prof. (Dr.) Vinod Kumar**

Department of Mathematics, Statistics & Computer Science,  
G. B. Pant University of Agriculture & Technology, Pantnagar, INDIA.

Research Areas - *Mathematical Statics, Time Series, Reliability Theory.*

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Department of Mathematics,  
Graphic Era Hill University, Dehradun Campus, INDIA.

Research Areas - *Operations Research, Reliability Theory, Fuzzy Reliability, System Engineering.*

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Faculty of Science and Engineering,  
Department of Mechanical Engineering,  
Université Laval, Quebec, CANADA.

Research Areas - *Industrial Engineering & Operations Research, Design & Planning of Supply Chains, Reliability & Maintenance, Quality, Decision Making, Modeling & Simulation, Stochastic Optimization, Game Theory.*

**Dr. Arpan Gupta**

School of Engineering,  
Indian Institute of Technology, Mandi, INDIA.

Research Areas - *Acoustics, Metamaterials, Vibrations, Bio-Mechanics, Computational Methods, Robotics.*

**Dr. Rajesh Joshi**

G. B. Pant Institute of Himalayan Environment and Development,  
An autonomus Institue of Ministry of Environment, Forest & Climate Change, Government of India,  
Kosi-Katarmal, Almora, INDIA.  
Research Areas - *Soft Computing Techniques, Environmental & Hydrological Modeling, Climate Impact Studies.*

**Prof. (Dr.) Alex Karagrigoriou**

Department of Mathematics,  
Division of Statistics & Actuarial-Financial Mathematics,  
University of the Aegean, Karlovasi, GREECE.  
Research Areas - *Model Selection, Information Theory & Divergence Measures, Goodness of Fit Tests, Reliability Theory, Applied Probability.*

**Dr. Anuj Kumar**

Department of Mathematics,  
University of Petroleum & Energy Studies, Dehradun, INDIA.  
Research Areas - *Reliability Theory, Wavelet Theory, Nature Inspired Optimization Algorithms.*

**Prof. (Dr.) Naoto Kaio**

Department of Economic Informatics, Faculty of Economic Sciences,  
Hiroshima Shudo University, Hiroshima, JAPAN.  
Research Areas - *Mathematical Theory of Reliability, Maintainability.*

**Dr. Lokesh Kumar Joshi**

Department of Mathematics,  
Faculty of Engineering and Technology,  
Gurukula Kangari University, Haridwar, INDIA.  
Research Areas - *Fuzzy Logic Applications.*

**Dr. Dinesh Bisht**

Department of Mathematics,  
Jaypee Institute of Information Technology, Noida, INDIA.  
Research Areas - *Soft Computing, Particle Swarm Optimization.*

**Dr. Navneet Joshi**

Department of Mathematics,  
Graphic Era Hill University, Bhimtal Campus, Nainital, INDIA.  
Research Areas - *Computational Fluid Dynamics, Mathematical Modelling, Simulation, Heat & Mass Transfer.*

**Dr. William Vesely**

Technical Risk Assessment,  
Office of Safety and Mission Assurance,  
NASA Headquarters, Washington, USA.  
Research Areas - *Aerospace Engineering, Fault Detection, Reliability Engineering, Risk Assessment.*

**Dr. Yuriy V. Kostyuchenko**

Workgroup on Disaster Study and Risk Analysis,  
Scientific Centre for Aerospace Research of the Earth,  
National Academy of Sciences of Ukraine, Honchar Street, Kiev, UKRAINE.  
Research Areas - *Risk Assessment, Environmental Studies, Vulnerability and Reliability.*

**Dr. Phuc Do**

Lorraine University,  
Campus Sciences BP09, 06 Vandoeuvre Cedex, FRANCE.  
Research Areas - *Reliability, Maintenance, Prognostics.*

**Prof. (Dr.) Bharatendra K. Rai**

Department of Decision and Information Sciences,  
Charlton College of Business,  
University of Massachusetts – Dartmouth, USA.  
Research Areas - *Multivariate Diagnosis, Pattern Recognition & Data Mining, Developing Meta-Models Using Computer Experiments, Prediction of Unexpended Warranty Costs, Field Performance Studies from Large Warranty Datasets.*

**Prof. (Dr.) John Andrews**

The Lloyd's Register Foundation Centre for Risk and Reliability Engineering,  
Nottingham Transportation Engineering Centre, Faculty of Engineering,  
University of Nottingham, Nottingham, UK.  
Research Areas - *System Reliability, Maintenance Modelling, Asset Management, Fault Diagnostics, Optimisation, Risk Assessment.*

**Prof. (Dr.) Tadashi Dohi**

Department of Information Engineering,  
Graduate School of Engineering,  
Hiroshima University, Hiroshima, JAPAN.  
Research Areas - *Reliability & Maintenance, Software Engineering, Dependable Computing.*

**Prof. (Dr.) Madhu Jain**

Department of Mathematics,  
Indian Institute of Technology Roorkee, Roorkee, INDIA.  
Research Area - *Stochastic Processes, Queuing Theory, Inventory Models.*



**Dr. Sangeeta Pant**

Department of Mathematics,  
University of Petroleum & Energy Studies, Dehradun, INDIA.  
Research Area - *Nature Inspired Optimization Techniques, Reliability Optimization.*

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Department of Mathematics,  
University of Petroleum & Energy Studies, Dehradun, INDIA.  
Research Area - *Fuzzy Reliability, Reliability Theory.*

**Prof. (Dr.) Revaz Kakubava**

Scientific-Research Center for System Problems,  
Georgian Technical University, Kostava str., Tbilisi, GEORGIA.  
Research Areas - *Reliability Theory, Queuing Theory, Mathematical Modeling & Optimization of Complex Systems.*

**Prof. (Dr.) S. C. Dimri**

Department of Computer Science and Engineering,  
Graphic Era (Deemed to be University), Dehradun, INDIA.  
Research Areas - *Advance Algorithms, Computer Graphics, Network Optimization.*

**Prof. (Dr.) Priti Kumar Roy**

Centre for Mathematical Biology and Ecology,  
Department of Mathematics,  
Jadavpur University, Kolkata, INDIA.  
Research Area - *Mathematical Biology on Ecology with Special Emphasis on Transmissible Disease, Modeling on Disease Psoriasis, HIV, Enzyme Kinetics, Bio-Diesel Production.*

**Prof. (Dr.) Lirong Cui**

School of Management and Economics,  
Beijing Institute of Technology, Beijing, CHINA.  
Research Areas - *Stochastic Modeling, Quality & Reliability Engineering, Simulation & Optimization, Operations Research, Applications of Probability & Statistics.*

**Prof. (Dr.) Won Young Yun**

Department of Industrial Engineering,  
Pusan National University, Pusan, REPUBLIC OF KOREA.  
Research Areas - *System Reliability Analysis, Maintenance Optimization, Warranty Policies.*

**Prof. (Dr.) Yi-Kuei Lin**

Department of Industrial Engineering and Management,  
Department of Electrical and Computer Engineering,  
National Chiao Tung University,  
1001 University Rd., Hsinchu 30010, TAIWAN.

Research Areas - *Network Analysis, System Reliability, Performance Evaluation, Service Management, Information Ethics.*

**Prof. (Dr.) Gopi Chattopadhyay**

Faculty of Science and Technology,  
Federation University, Mount Helen VIC 0, AUSTRALIA.

Research Areas - *Asset Management, Maintenance & Reliability Engineering, Operations Management, Project Management & Quality Engineering.*

**Dr. Adarsh Anand**

Department of Operational Research,  
University of Delhi, New Delhi, INDIA.

Research Areas - *Innovation Diffusion & Multi Generation Modeling in Marketing, Software Reliability Growth Modeling.*

**Prof. (Dr.) H. M. Srivastava**

Department of Mathematics and Statistics,  
University of Victoria, Victoria, British Columbia VW R, CANADA.

Research Areas - *Analytic Number Theory, Inventory Modelling & Optimization, Fractional Calculus, q-Series & q-Polynomials, Real & Complex Analysis, Fractional Calculus, Integral Equations & Transforms, Higher Transcendental Functions.*

**Prof. (Dr.) Patricio Franco**

Department of Materials and Manufacturing Engineering,  
Technical University of Cartagena, Murcia, SPAIN.

Research Areas - *Manufacturing Processes, Dimensional Metrology, Metal Forming Processes, Joining & Welding Processes, Optimization Of Manufacturing Systems, Numerical Modelling Of Manufacturing Processes, CAD/CAE/CAM/CAPP, Quality Assurance, Productivity Improvement.*

**Prof. (Dr.) Suresh Kumar Sharma**

Centre for Systems Biology and Bioinformatics,  
Panjab University, Chandigarh, INDIA.

Research Areas - *Statistical Modelling, Biostatistics, Ranking & Selection, Statistical Inference, Distribution Theory.*

**Dr. Anatoly Lisnianski**

The Israel Electric Corporation Ltd.,  
Planning, Development & Technology Division, Haifa, ISRAEL.

Research Areas - *Reliability, Risk Analysis, Maintenance.*

**Dr. Sanjay Kumar Chaturvedi**

Subir Chowdhury School of Quality and Reliability,  
Indian Institute of Technology Kharagpur, Kharagpur, INDIA.

Research Areas - *System Reliability Modelling & Analysis, Reliability Data Analysis, Reliability Estimation, Maintenance Engineering.*

**Dr. Geeta Arora**

Department of Mathematics,  
Lovely Professional University, Phagwara, INDIA.

Research Areas - *Numerical Analysis, Partial Differential Equations.*

**Prof. (Dr.) J. K. Pandey**

Department of Research & Development,  
University of Petroleum & Energy Studies, Dehradun, INDIA.

Research Areas - *Nano-technology, Bio-polymers, Water Treatment and Composites.*

**Dr. Jitendra Pal Singh**

Advanced Analysis Center,  
Korea Institute of Science and Technology, Seoul, SOUTH KOREA.

Research Areas - *Nanomaterials, Magnetism, Electronic Structure.*

**Dr. Surya Parkash**

National Institute of Disaster Management (NIDM),  
Ministry of Home Affairs, Government of India, New Delhi, INDIA.

Research Areas - *Landslides, Avalanches, Earthquakes, Floods, Tsunami, GLOFs, LLOFs, Dam Bursts, Mining Disasters, Community Based Disaster Risk Management, Damage and Loss Assessment, Disaster Safe Hill Area Development, Environmental Management.*

**Dr. Sachin Kumar Mangla**

Plymouth Business School,  
University of Plymouth, UK.

Research Areas - *Green Supply Chain Management, Sustainability, Risk Management, Decision Support System/MCDM, Operations Management.*

**Dr. Rajeev Singh**

Department of Computer Engineering,  
G. B. Pant University of Agriculture & Technology, Pantnagar, INDIA.

Research Areas - *Data Structures, Algorithms, Computer Networks, Network Security.*

**Dr. Santosh Kumar**

Department of Computer Science and Engineering,  
Graphic Era (Deemed to be University), Dehradun, INDIA.

Research Areas - *Ubiquitous Computing, Network Security, Internet on things, Software Engineering.*

**Dr. Amit Kumar**

Department of Mathematics,  
Lovely Professional University, Phagwara, INDIA.  
Research Areas - *Reliability Theory, System Engineering, Multi-State Systems.*

**Dr. Nupur Goyal**

Department of Mathematics,  
Roorkee Institute of Technology, Roorkee, INDIA.  
Research Areas - *Reliability Theory, System Engineering, Multi-State Systems.*

**Prof. (Dr.) Ali Muhammad Ali Rushdi**

Department of Electrical and Computer Engineering,  
King Abdulaziz University, Jeddah 19, SAUDI ARABIA.  
Research Areas - *System Reliability, Boolean Analysis and Reasoning, Analysis and Design of Algorithms, Engineering Mathematics.*

**Prof. (Dr.) G. S. Ladde**

Department of Mathematics and Statistics,  
University of South Florida, 0 East Fowler Avenue, Tampa, Florida, USA.  
Research Areas - *Hereditary Dynamic Systems; Local Lagged Generalized Method of Moments (LLGMM); Multi-agent & Multi-Market/Finance; Multivariate/Large-Scale Systems Analysis; Stability Stochastic Hybrid Dynamical; Stochastic Modeling of Dynamical Processes in Biological, Chemical, Engineering, Medical, Physical and Social Sciences; Stochastic Modeling of Network Dynamic; Stochastic System Parameter and State Estimation; Time Series Analysis and Applications.*

**Dr. Harish Garg**

School of Mathematics,  
Thapar Institute of Engineering and Technology University, Patiala, INDIA.  
Research Areas - *Multi Criteria Decision Making, Reliability, Fuzzy Optimization, Multi Objective Optimization, Aggregation Operators, Computational Intelligence, Optimization Techniques, Various Nature Inspired Algorithms (E.G. Genetic Algorithms, Swarm Optimization, Artificial Bee Colony Etc.), Fuzzy and Intuitionistic Fuzzy Set Theory, Expert Systems. Application Areas Include Wide Range of Industrial and Structural Engineering Design Problems.*

**Prof. (Dr.) Santosh Kumar**

Department of Mathematics and Statistics,  
University of Melbourne, Melbourne, AUSTRALIA.  
Research Areas - *Operations Research, Optimization Techniques, Applied Mathematics.*

**Prof. (Dr.) Ljubisa Papic**

DQM Research Center, Prijevor, SERBIA.

Research Areas - *Reliability Engineering, Maintainability Engineering, Failure Analysis, Safety Analysis, Quality Management.*

**Prof. (Dr.) Piotr Kulczycki**

Centre of Information Technology for Data Analysis Methods,

Systems Research Institute,

Polish Academy of Sciences, Warsaw, POLAND.

Research Areas - *Data Analysis, Mathematical Statistics, Control Engineering, Mathematical Modelling, Computational Intelligence.*

**Dr. Garima Chopra**

Department of Mathematics,

University Institute of Engineering & Technology,

Maharshi Dayanand University, Rohtak, INDIA.

Research Areas - *Wavelet Analysis, Image Processing, Reliability Analysis.*

**Prof. (Dr.) P. K. Kapur**

Former Professor, University of Delhi,

Amity Center for Interdisciplinary Research,

Amity University, Noida, INDIA.

Research Areas - *Software Reliability, Modeling, Optimization, Marketing Research, Vulnerability Discovery Modeling.*

**Dr. Ritu Arora**

Department of Mathematics,

Gurukul Kangari University, Haridwar, INDIA.

Research Areas - *Mathematical Modeling, Fuzzy Logic, Wavelet Theory, Applied Mathematics.*

**Prof. (Dr.) Chandra K. Jaggi**

Department of Operational Research,

Faculty of Mathematical Sciences,

New Academic Block, University Of Delhi, Delhi, INDIA.

Research Areas - *Supply Chain Management, Inventory Management, Trade Credit, Two Warehousing Models.*

**Prof. (Dr.) J. P. Singh Joorel**

Department of Statistics,

University of Jammu, Jammu, INDIA.

Research Areas - *Operations Research, Reliability Theory, Sampling & amp; Statistical Inference.*

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Department of Technology Systems,

College of Engineering and Technology,  
East Carolina University, USA.

Research Areas - *Operations Management, Inventory Planning, Supply Chain Management, Operations Research, Management Science; Quality Assurance.*

**Prof. (Dr.) Jianqiang Gao**

College of Computer and Information Engineering,  
Hohai University, Nanjing 1009, CHINA.

Research Areas - *Remote Sensing Image Classification, Image Processing, Information Processing System and Pattern Recognition.*

**Dr. Bidyut Bikash Gogoi**

Space Navigation Group,  
Indian Space Research Organization, ISRO, INDIA.

Research Areas - *u1-69Heat and Mass Transfer, Mathematical Biology, Computational Fluid Dynamics, Numerical Methods for PDEs, Scientific Computing, Gas Dynamics, Lattice Boltzmann Methods, Time Synchronization, Navigation, Precision Farming, Numerical Analysis.*

**Dr. Bhagawati Prasad Joshi**

Department of Applied Sciences,  
Seemant Institute of Technology, Pithoragarh, INDIA.

Research Areas - *Soft Computing Techniques: Theories and Applications, Reliability Theory, Decision Making under Uncertain Environment.*

**Prof. (Dr.) Hoang Pham**

Department of Industrial & Systems Engineering,  
School of Engineering,  
Rutgers State University, New Jersey, USA.

Research Areas - *Reliability Engineering, Optimization, Statistical Inference.*

**Prof. (Dr.) Antonella Petrillo**

Department of Engineering,  
University of Napoli "Parthenope", ITALY.

Research Areas - *Operation Research, Industrial Engineering, Product Life Cycle Management, Business Model, Decision Analysis and Decision Support System, Modeling and Simulation, Safety-Critical System Performance.*

**Prof. Xiao-Zhi Gao**

School of Computing,  
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Research Areas - *Artificial Intelligence, Machine Learning, Data Mining, Optimization.*

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Faculty of Transport,  
Transport Systems and Traffic Engineering Department,  
Silesian University of Technology, Gliwice, POLAND.  
Research Areas - *Traffic Engineering, Traffic Modelling, Driver Behaviour Analysis, Transport Statistics, Optimization in Transport.*

**Prof. (Dr.) Sanjeev Kumar Singh**

Department of Mathematics,  
University of Petroleum & Energy Studies, Dehradun, INDIA.  
Research Areas - *Differential Geometry, Mathematical Modeling, Soft Computing, Data Mining.*

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Director of the Monitoring Centre for Science and Education,  
Department of Economics and Management in Machine Building,  
Saint Petersburg Polytechnic University, Saint Petersburg, RUSSIA.  
Research Areas – *Statistical Quality Control and Quality Management, Systems Reliability.*

**Dr. A. P. Singh**

Department of Mathematics,  
S. G. R. R. (P. G.) College, Dehradun,  
H.N.B. Garhwal Central University, Srinagar (Garhwal), INDIA  
Research Areas - *Operations Research, Inventory Management, Optimization.*

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Computer Science Faculty, Core Competence Center C for Operations Research,  
Chair for Operations Research,  
Universität der Bundeswehr München, Neubiberg, GERMANY.  
Research Areas - *Optimization of Complex Systems, IT Based Decision Support Systems/ Reachback Architectures, Strategic Management.*

**Prof. (Dr.) Alexander Bochkov**

Sustainability Analysis in the Oil and Gas Industry Department,  
Centre "Risk Analysis", LLC NIIGAZECONOMIKA, Moscow, RUSSIA.  
Research Areas - *Non-Stationary Processes, Risk Assessment, Analysis and Management, System Analysis and Operation Research, Vulnerability and Survivability of Large-Scale Systems.*

**Prof. (Dr.) Anu A. Gokhale**

Department of Technology,  
Illinois State University, Normal, USA.  
Research Areas - *Analytics, Computational Algorithms, Data Mining.*

**Dr. Sunil Luthra**

Department of Mechanical Engineering,  
State Institute of Engineering & Technology (Formerly known as Government Engineering College) Nilokheri, Haryana, INDIA  
Research Areas - *Operations Management, Green/Sustainable Supply Chain Management, Production and Industrial Engineering, Renewable/Sustainable Energies.*

**Prof. (Dr.) Liudong Xing**

Department of Electrical and Computer Engineering,  
University of Massachusetts, Dartmouth, USA.  
Research Areas - *Systems Reliability Engineering.*

**Prof. (Dr.) Sankar Kumar Roy**

Department of Applied Mathematics with Oceanology and Computer Programming,  
Vidyasagar University, Midnapore, INDIA.  
Research Areas - *Operations Research- Transportation Problem, Game Theory, Inventory Management and Facility Location Problem.*

**Prof. (Dr.) Paulo Peças**

IDMEC, Instituto Superior Técnico,  
Universidade de Lisboa, Lisboa, PORTUGAL.  
Research Areas - *Life Cycle Engineering, Life Cycle Cost Modelling, Life Cycle Analysis, Eco-Efficiency, Sustainable Production, Lean Manufacturing, Lean Digitalization, Lean & Industry 4.0 (Lean 4.0), Additive Manufacturing, Natural Fibre Composites.*

**Dr. Vijay Kumar**

Department of Mathematics,  
AIAS, Amity University, Noida, INDIA.  
Research Areas - *Reliability Engineering, Optimization, Modeling, Optimal Control Theory.*

**Dr. Deepti Aggrawal**

University School of Management & Entrepreneurship,  
Delhi Technological University, East Delhi Campus, INDIA.  
Research Areas - *Operational Research, Decision Sciences, Marketing Analytics.*

**Dr. Saurabh Kapoor**

Department of Education in Science and Mathematics,  
Regional Institute of Education, Bhubneshwar,  
(National Council of Educational Research and Training, Delhi),  
Ministry of HRD, Government of India, INDIA.  
Research Areas - *Computational Fluid Dynamics, Hydrodynamics Stability of Flow through Porous Medium, Applied Numerical Method, FDM, FEM, SCM, B Spline FEM, ICT in Mathematics, Mathematics Education.*



**Prof. (Dr.) Nita H. Shah**

Department of Mathematics,  
University School of Sciences, Gujarat University,  
Ahmedabad – 380 009, Gujarat, INDIA.  
Research Areas - *Operations Research, Dynamical Systems, Applied Mathematics.*

**Prof. (Dr.) Ajit Kumar Verma**

Department of Fire Safety and HSE Engineering,  
Western Norway University of Applied Sciences,  
Haugesund, NORWAY.  
Research Areas - *Reliability in Engineering Design, Software Reliability, Design and Analysis of Fault Tolerant Systems, Physics of Failures, Power System Reliability.*

**Dr. Chandra Shekhar**

Department of Mathematics,  
Birla Institute of Technology & Science,  
Pilani, 333031, Rajasthan, INDIA.  
Research Areas - *Queueing theory, Stochastic processes, Computer and communication system, Reliability and maintainability, Manufacturing and machine repair problem, Reliability engineering, Fuzzy set and Fuzzy logic, Inventory Theory, Statistical inference and analysis, Evolutionary and Nature Inspired Optimization Techniques.*

**Dr. Mukesh Kumar**

Department of Mathematics,  
Motilal Nehru National Institute of Technology Allahabad, Prayagraj, INDIA.  
Research Areas - *Nonlinear Partial Differential Equations, Similarity Transformation Method, Lie Group Theory.*

**Prof. (Dr.) Yigit Kazancoglu**

Department of International Logistics Management,  
Yasar University, Izmir, TURKEY.  
Research Areas - *Operations Management, Supply Chain Management, Multi-Criteria Decision Making, Total Quality Management.*

**Dr. Shailendra Rajput**

Department of Electrical and Electronic Engineering,  
Ariel University, P.O.B. 3, Ariel, ISRAEL.  
Research Areas - *Energy Harvesting, Dielectric Materials, Piezoelectricity and Ferroelectricity.*

**Prof. (Dr.) Xufeng Zhao**

Nanjing University of Aeronautics and Astronautics,

No. 29, Jiangjun Avenue, Nanjing 211106, CHINA.

Research Areas - *Quality Management, Reliability Engineering, Maintenance Modeling and Optimization, Management of Information and Industrial Systems.*



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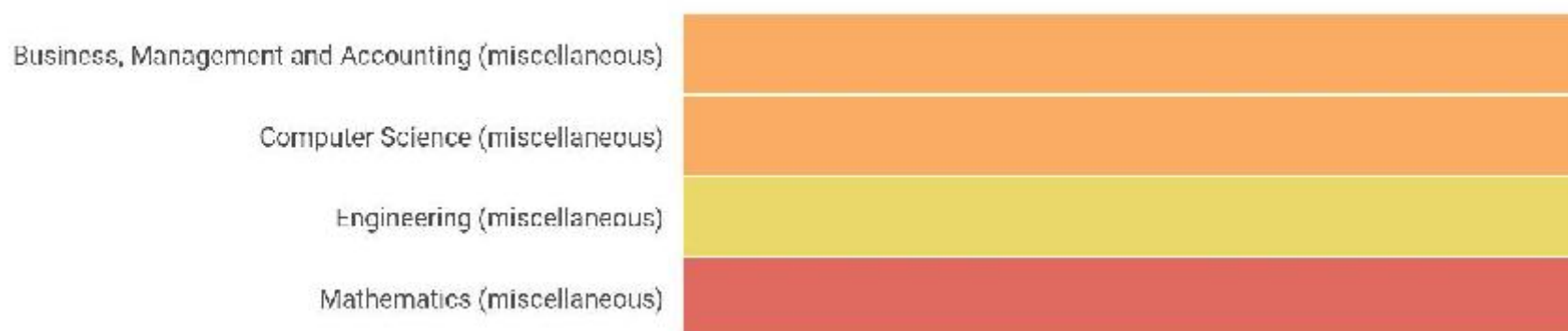
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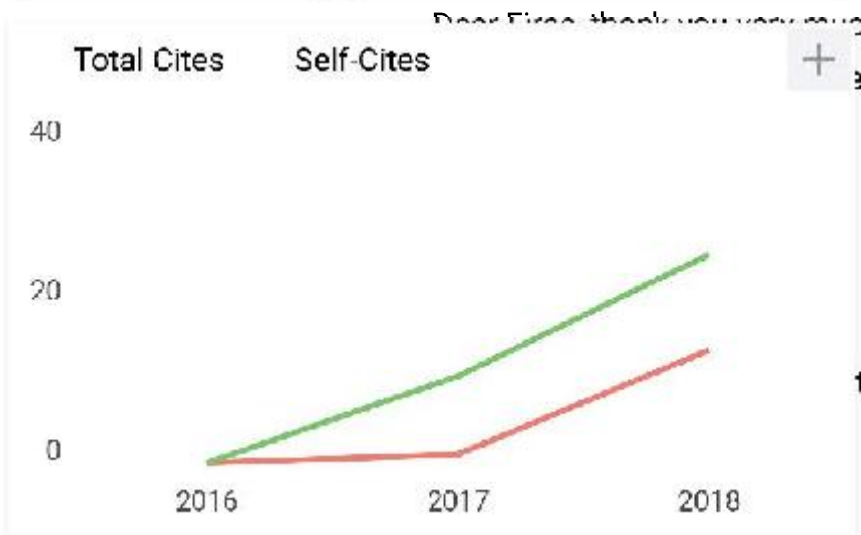
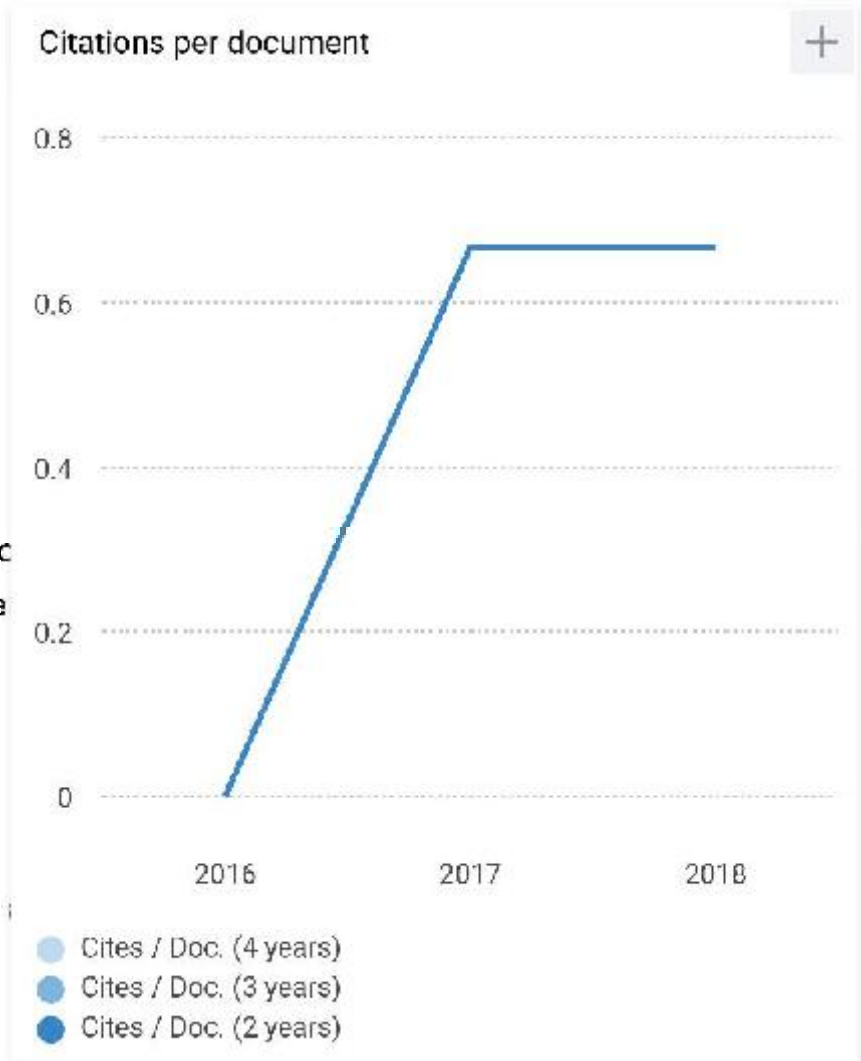
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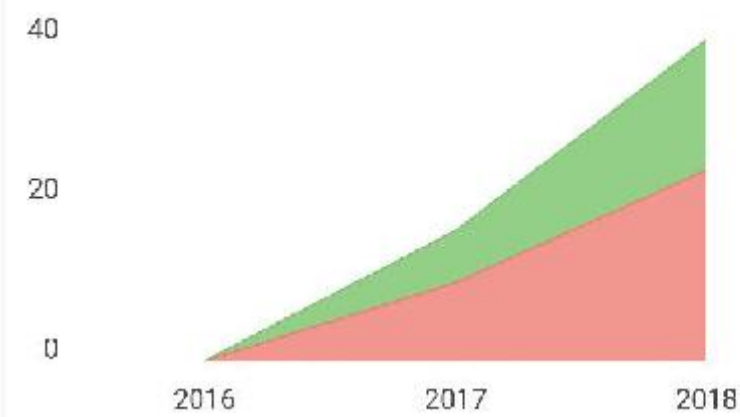
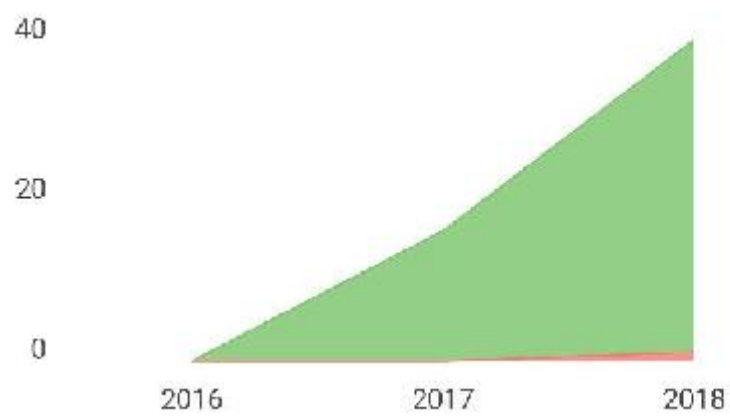
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## Dynamics of A Re-Parametrization of A 2-Dimensional Mapping Derived from Double Discrete Sine-Gordon Mapping

**La Zakaria**

Department of Mathematics,  
Faculty of Mathematics and Natural Sciences, Universitas Lampung, Indonesia.  
*Corresponding author:* lazakaria.1969@fmipa.unila.ac.id

**Johan Matheus Tuwankotta**

Analysis and Geometry Group,  
Faculty of Mathematics and Natural Sciences, Institut Teknologi Bandung, Indonesia.  
E-mail: theo@dns.math.itb.ac.id

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### Abstract

We study the dynamics of a two dimensional map which is derived from another two dimensional map by re-parametrizing the parameter in the system. It is shown that some of the properties of the original map can be preserved by the choice of the re-parametrization. By means of performing stability analysis to the critical points, and also studying the level set of the integrals, we study the dynamics of the re-parametrized map. Furthermore, we present preliminary results on the existence of a set where iteration starts at a point in that set, in which it will go off to infinity after finite step.

**Keywords-** Re-parametrizing, 2-dimensional mapping, Generalized double discrete sine-Gordon, Integral.

### 1. Introduction

Arguably, one of the most important and general integrable maps is known in the literature as the Quispel-Roberts-Thompson map (QRT). It is a two-dimensional map depending on 18 parameters. The QRT map is closely related to so called soliton equations (Quispel et al., 1988; Quispel et al., 1989). More recent studies have focused on generalizations of QRT maps. One of them was proposed by Joshi and Kassotakis (2019). Main result of their paper is a new connection between two major theories that generalize QRT maps. They provide a new formulation of QRT involutions in terms of Hirota derivatives and discover conditions under which each involution can be factorized into two further involutions.

The sine-Gordon equation is a partial differential equation, which is known to have soliton solutions; hence, it is a soliton equation (Quispel et al., 1991). Discretizations of the sine-Gordon equation have been done in various ways (Quispel et al., 1988; Quispel et al., 1989; Quispel et al., 1991). The reduction of the sine-Gordon equation to a two dimensional ordinary difference equations using a standard staircase (see Van der Kamp and Quispel, 2010 for the method) is known as being a special case of the celebrated QRT map. Recently, Celledoni et al. (2019) have studied a new systematic approach for calculating the preserved measures and integrals of a rational map in which a two-dimensional sine-Gordon (standard) map was chosen as an example (see Celledoni et al., 2019 §4.3 Example 4).

Generating a new family of mapping from a known one is an interesting topic to study. A piece-wise linear map from a known integrable map by using the ultra discretization technique is

generated by Tuwankotta et al. (2004). The number of independent integrals are preserved by the transformation which implies that the integrability is preserved. In our article (Zakaria and Tuwankotta, 2016) a straight-forward generalization by adding parameters in the Lax pair of the ordinary discrete sine-Gordon partial difference equation has done. By using the standard staircase method, the resulted equation is then reduced to system of ordinary difference equations (see Van der Kamp et al., 2007 for the method). Note that this generalized sine-Gordon system is also analyzed in (Duistermaat, 2010).

Another novel method to introduce a new class of discrete systems from an integrable discrete system, is by introducing the concept of dual (Quispel et al., 2005). This works beautifully for a single discrete equation, although the resulting equation may not be new nor integrable. This idea of dual is extended to system of discrete equations in (Tuwankotta et al., 2019). The latter is interesting in the sense that the method proposed there produces in general more than one system.

Roberts et al. (2002) constructed a new family of mapping by interchanging the parameter and the integral. For example, consider a discrete dynamical system in  $\mathbb{R}^n$  which is denoted by:  $\mathbf{x} = \mathbf{f}(\mathbf{x}, \lambda)$ , where the prime denotes the upshift,  $\lambda$  is a parameter in the system, and  $\mathbf{f} : \mathbb{R}^n \times \mathbb{R} \rightarrow \mathbb{R}^n$ . We assume that there exists a smooth function  $G : \mathbb{R}^n \times \mathbb{R} \rightarrow \mathbb{R}$  such that:  $G(\mathbf{x}, \lambda) = G(\mathbf{x}', \lambda')$ , which is called an integral for the system. Suppose that we can solve the equation  $G(\mathbf{x}, \lambda) = 0$  for  $\lambda = \lambda(\mathbf{x})$ . Then by substituting this solution to  $\mathbf{f}(\mathbf{x}, \lambda)$  (and call it  $g(\mathbf{x})$ ), we derive a new discrete system:  $\mathbf{x}' = g(\mathbf{x})$ , with integral:  $G(\mathbf{x}, \lambda(\mathbf{x}))$ .

We will follow this technique and apply it to a generalized sine-Gordon equation. The aim of this paper is to show a number of properties of the new mapping (after re-parameterization) and to compare them with the original mapping. Furthermore, the qualitative behavior of the new mapping is studied by means of obtaining fixed points and their stability, and also the base points.

The outline of this paper is the following. In Section 2, a system of first order difference equations derived from generalized double discrete sine-Gordon (double sine-Gordon) equations is formulated by restricting to traveling wave solution. In Section 3 a new integrable mapping derived from sine-Gordon mapping by interchanging the role of the integral and the parameter in the original system is presented. The new system is then analyzed by means of describing its symmetry and measure preservation. In Section 4 finding fixed points and periodic points (and their linear stability) and also computing the base points are discussed. These are done in the remaining subsections in Section 4. We have divided Section 4 into five subsections. This paper ends with some concluding remarks in Section 5.

## 2. Formulation of the Problem

Consider a three parameters family of partial difference equation on two dimensional lattice:

$$V_{l,m} V_{l+1,m} V_{l,m+1} V_{l+1,m+1} = \lambda V_{l,m} V_{l+1,m} V_{l,m+1} V_{l+1,m+1} \quad (1)$$

This equation is derived from the compatibility condition of the generalized Lax pair of the classical discrete sine-Gordon equation (see Zakaria and Tuwankotta, 2016). The travelling wave solutions of (1) can be obtained by considering the following form:

$$V_{l,m} = V_n, n = z_1 l + z_2 m,$$

where  $z_1$  and  $z_2$  are relatively prime integers. Substituting this to (1), we have an infinite hierarchy of mapping labeled by  $z_1$  and  $z_2$ , i.e.

$$V_{n-1} = \frac{V_n}{z_1}, V_{n-2} = \frac{V_n}{z_1^2}, \dots, V_{n-3} = \frac{V_n}{z_1^3}, \dots \quad (2)$$

If  $z_1$  and  $z_2$  are fixed, a mapping from  $\mathbb{R}^{z_1+z_2}$  to  $\mathbb{R}^{z_1+z_2}$  can be obtained from (2). It can be noted that by setting  $z_1 = 2$  and  $z_2 = 1$  in (2), we can derive the two dimensional mappings in Quispel et al., (1991).

Let  $z_1 = z_2 = 1$ ,  $z_1 = 2$ , and  $z_2 = 1$ . From (2), we derive:

$$V_{n-2} = \frac{V_n^2}{V_n + V_{n-1}^2},$$

which is a second order difference equation. Then, by writing  $x_n = V_{n-1}$ , and  $y_n = V_n$ .

We derive a system of first order difference equations:

$$\begin{aligned} x_{n-1} &= \frac{x_n^2}{y_n + x_n^2}, \\ y_{n-1} &= x_n. \end{aligned} \quad (3)$$

Let us define:  $f : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ , by

$$f(x, y) = \left( \frac{x^2}{y + x^2}, x \right).$$

Using this, we can write (3) as:

$$(x, y) = f(x, y) \quad (4)$$

where the prime denotes the upshift. The mapping (4) has an integral, i.e.

$$F(x, y) = \frac{x}{y} - \frac{y}{x} - xy - \frac{1}{xy}. \quad (5)$$

Thus,  $F(x_{n-1}, y_{n-1}) = F(x_n, y_n)$  for all  $n \in \mathbb{N}$ .

### 3. Reparametrized Mapping and Its Properties

Consider the mapping in (4). For  $z_1 = 1$ , we solve the equation  $F(x, y) = 0$  for  $y$  to derive:

$$(x, y) = \frac{1 - x^2 y^2}{x^2 - y^2}. \quad (6)$$

Substituting (6) into (4), we derive a new mapping:

$$(x, y) = \hat{f}(x, y) = (y, x), \quad (7)$$

with integral

$$\hat{F}(x, y) = \frac{1 - x^2 y^2}{x^2 - y^2}. \quad (8)$$

Note that the integral (8) is nothing but the function  $(x, y)$  in (6). The new map,  $\hat{f}(x, y)$ , has the following properties:

The orbits of  $\hat{f}(x, y)$  is 4-periodic. This is simple to show that the linear map in (7) is trivial, its second iterate leads to  $(x, y) = (x, y)$  thus all the results are immediate (the fourth iterate is the identity function).

$\hat{f}(x, y)$  is area preserving. A two-dimensional map is area preserving (also called conservative) when its Jacobian determinant  $|J(x, y)|$  is equal to 1 in all the points  $(x, y)$  of the plane, which for the linear map in equation (7) is immediate (the Jacobian determinant is constant and equal to 1), as well as the so-called "reverse symmetry", since  $\hat{f}^{-1}(u, v) = \hat{f}(u, v)$ . Alternatively, we can use different procedure to show that an integrable map is measure preserving (Roberts et al., 2002). For our map,  $\hat{f}(x, y)$ , the procedure can be followed as shown below.

The mapping  $\hat{f}(x, y)$  is measure preserving (area preserving) because there is the density  $\hat{\rho}(x, y)$  such that

$$|D\hat{f}(x, y)| = \frac{\hat{\rho}(x, y)}{\hat{\rho}(x, y)} = 1$$

when the density  $\hat{\rho}$  is given by

$$\hat{\rho}(x, y) = \frac{1}{xy} \frac{F(x, y)}{x^2 - y^2} = \frac{1}{x^2 - y^2}.$$

Consider  $G_1(x, y) = (y, x)$ . Note that  $G_1(G_1(x, y)) = G_1(y, x) = (x, y)$ . This implies that:

$$G_1^{-1} = G_1. \text{ Since } G_1 \circ \tilde{f} \circ G_1^{-1} = \tilde{f}^{-1} \text{ then } G_1 \text{ is a reversing symmetry for } \hat{f}.$$

There exists a symmetry  $S_1(x, y) = (x, -y)$  such that  $S_1 \circ \hat{f} \circ S_1^{-1} = \hat{f}$ .

The dynamics of the mappings in (7) for  $\hat{F} = 1$  on every level set  $\hat{F} = c$  is basically identical to the dynamics of the mappings in (4) on the level set  $F = 0$  for  $\hat{F} = c$ . Furthermore, this provides

us with the existence of a 4-periodic points of the mappings: (7) for  $\mu = 1$  for every  $\mu$ . The locations of these points are in the level set  $F = 0$ .

Still fixing the value of  $\mu = 1$ , let us now, reparametrize the parameter in (4) by  $a_0, a_1$ . It follows immediately that the map is given by

$$\tilde{f}(x, y) = \frac{1}{y} \frac{a_0 + a_1 x^2}{x^2 + a_0 + a_1}, x. \quad (9)$$

Consequently, we transform the integral (5) to

$$\tilde{F}(x, y) = a_0 + a_1 \frac{x}{y} \frac{y}{x} - xy \frac{1}{xy}. \quad (10)$$

For the case where:  $a_1 \neq 0$ , let us add a constant in (10) which takes a special form:  $b_0, b_1$ , i.e.

$$\tilde{\tilde{F}}(x, y) = a_0 + a_1 \frac{x}{y} \frac{y}{x} - xy \frac{1}{xy} + b_0 + b_1, \quad (11)$$

Note that  $\tilde{\tilde{F}}(x, y)$  is linear in  $x$ . Furthermore:

$$\tilde{\tilde{F}} \circ \tilde{f} = \tilde{\tilde{F}} \left( \frac{1}{y} \frac{a_0 + a_1 x^2}{x^2 + a_0 + a_1}, x \right) = \tilde{\tilde{F}}(x, y).$$

This implies that we can look at the zero level curve of  $\tilde{\tilde{F}}$  and solve it for  $x$ :

$$\tilde{\tilde{f}}(x, y) = \frac{1}{x^2 a_1} \frac{x^2 y^2 - x^2 a_0 - y^2 a_0 - xy b_0}{y^2 a_1 - xy b_1}. \quad (12)$$

Substituting this expression of  $\tilde{\tilde{f}}(x, y)$  into  $\tilde{f}$  in (9) gives:

$$\hat{\tilde{f}}(x, y) = \frac{a_1 y - x^4 y - x^3 b_0 - x - x^2 a_0 - b_1}{a_1 - x^4 - xy b_0 - xy x^2 - a_0 - b_1}, x.$$

It is interesting to note that  $\tilde{\tilde{f}} \circ \hat{\tilde{f}}(x, y) = \tilde{\tilde{f}}(x, y)$  which implies that  $\hat{\tilde{f}}(x, y)$  is an integral for the system:

$$(x, y) = \hat{\tilde{f}}(x, y). \quad (13)$$

The mapping (13) has some properties:

$\hat{\tilde{f}}$  is measure preserving and orientation-reversing (or anti-measure preserving), which means (Roberts et al., 2002)

$$|D\hat{f}| = \frac{\hat{f}(x, y)}{\begin{vmatrix} a_1 & y & x^4 & y & x^3 b_0 & x & 1 & x^2 a_0 & b_1 \\ a_1 & 1 & x^4 & xy & b_0 & xy & x^2 & a_0 & b_1 \end{vmatrix}}, x$$

Note that from the right-hand side, the determinant of the Jacobian of  $\hat{f}$  is

$$\frac{b_1^2 x^2 \ a_0^2 x^2 \ a_0 \ x^4 \ 1 \ x^2 \ a_1 b_0 b_1 x^2 \ 2a_0 x^2 \ x^4 \ 1 \ a_1^2 \ b_0^2 \ 2 \ x^4 \ x^8 \ 1}{b_1 xy \ x^2 \ a_0 \ a_1 \ b_0 xy \ x^4 \ 1^2}$$

Meanwhile, from the right-hand side, we have

$$\frac{b_1^2 x^2 \ a_0 \ x^2 \ a_0 x^2 \ 1}{a_1 \ 1 \ x^4 \ xy b_0 \ xy \ x^2 \ a_0 \ b_1^2} \frac{a_1 b_0 b_1 x^2 \ 2a_0 x^2 \ x^4 \ 1}{a_1 \ 1 \ x^4 \ xy b_0 \ xy \ x^2 \ a_0 \ b_1^2}$$

$$\frac{a_1^2 \ b_0^2 \ 2 \ x^4 \ x^8 \ 1}{a_1 \ 1 \ x^4 \ xy b_0 \ xy \ x^2 \ a_0 \ b_1^2}$$

where the so-called density  $\hat{f}$  is given by

$$\hat{f}(x, y) = \frac{1}{x y} \tilde{F}^{-1} = \frac{1}{a_1(x^2 - y^2) b_1 x y}$$

The function  $\hat{G}_1(x, y) = y, x$ , is a reversing symmetry for  $\hat{f}$ .

#### 4. Critical Point and Base Point of the Integral

There are two important elements in analyzing the dynamics of system (13), i.e. Fixed Point (FP) and Base Point (BP). FP can be obtained by finding the critical point of the integral function, while BP is defined as the point where the integral function is singular. At the BP, level sets for various values of the integral function intersect each other.

##### 4.1 The Critical Point

The critical points of the integral function (12) are solutions of

$$\frac{\hat{f}}{x} = \frac{a_1 \ 2x \ y^4 \ 1 \ y \ x^2 \ y^2 \ b_0 \ y \ x^2 y^2 \ 1 \ x^2 \ y^2 \ a_0 \ b_1}{x^2 \ y^2 \ a_1 \ xy b_1^2} = 0,$$

$$\frac{\hat{f}}{y} = \frac{a_1 \ 2y \ x^4 \ 1 \ x^3 \ xy^2 \ b_0 \ x \ x^2 y^2 \ 1 \ x^2 \ y^2 \ a_0 \ b_1}{x^2 \ y^2 \ a_1 \ xy b_1^2} = 0.$$

To obtain the solutions, we can do as follow:

$$y \frac{\hat{\phantom{x}}}{x} - x \frac{\hat{\phantom{y}}}{y} = 0 \quad \frac{x^2 - y^2 - a_1 - 2xy - b_0 - x^2 y^2 - x^2 - y^2 - a_0 - b_1}{x^2 - y^2 - a_1 - xy b_1} = 0$$

From the left-hand side in the last equation, we have two lines,  $y = x$  and  $y = -x$  as solutions.

By substituting these into  $\frac{\hat{\phantom{x}}}{x} = 0$  or  $\frac{\hat{\phantom{y}}}{y} = 0$  and then solve it, we have

$$x, y = 1, -1 \text{ and } 1, \mp 1 \tag{14}$$

It is easy to verify that: (1,1) and (-1, -1) are two fixed points while (1, -1) and (-1, 1) are two 2-periodic points.

#### 4.2 The Base Point

Apart from the critical points, the so-called base points also play a crucial role in the dynamics of (13). Note that the invariant (12) can be written as a rational function. A point  $(x_0, y_0)$  is a base point if it is a common zero of the numerator and denominator of  $\hat{\phantom{x}}, \hat{\phantom{y}}$  equal to zero. In our case, we will discuss two conditions,  $a_1 = 0$  and  $a_1 \neq 0$ .

For  $a_1 = 0$ , the base points are

$$(x_0, y_0) = 0, \sqrt{\frac{1}{a_0}}, 0, -\sqrt{\frac{1}{a_0}}, \sqrt{\frac{1}{a_0}}, 0, -\sqrt{\frac{1}{a_0}}, 0 \tag{15}$$

Note that the points in eq. (15) are base points of the system (13) when the parameter  $a_1 = 0$ . Two

base points,  $(-\sqrt{\frac{1}{a_0}}, 0)$ , are directly mapped to infinity by the mapping  $\hat{\phantom{x}}, \hat{\phantom{y}}$ , in the sense that one

of both of the component of  $\hat{\phantom{x}}, \hat{\phantom{y}}$  blows up at these points. And then the other points,  $(0, \sqrt{\frac{1}{a_0}})$ ,

are base points in which they are mapped to other base points after one iterate of  $\hat{\phantom{x}}, \hat{\phantom{y}}$ .

For  $a_1 \neq 0$ , the base points can be obtained by solving  $1 - x^2 y^2 - x^2 a_0 - y^2 a_0 - xy b_0 = 0$  and  $x^2 a_1 - y^2 a_1 - xy b_1 = 0$ . Based on our computations, the base points are

$$(x_0, y_0) = (x_{10}, y_{10}), (\mp x_{10}, y_{10}), (x_{20}, y_{20}), (\mp x_{20}, y_{20})$$

where



$$x_{10} \sqrt{\frac{d^2 b_0 a_0 d^2 8 d d^2 2 a_0 d a_0 d^2 b_0 b_0 d}{2 d d^2 2}};$$

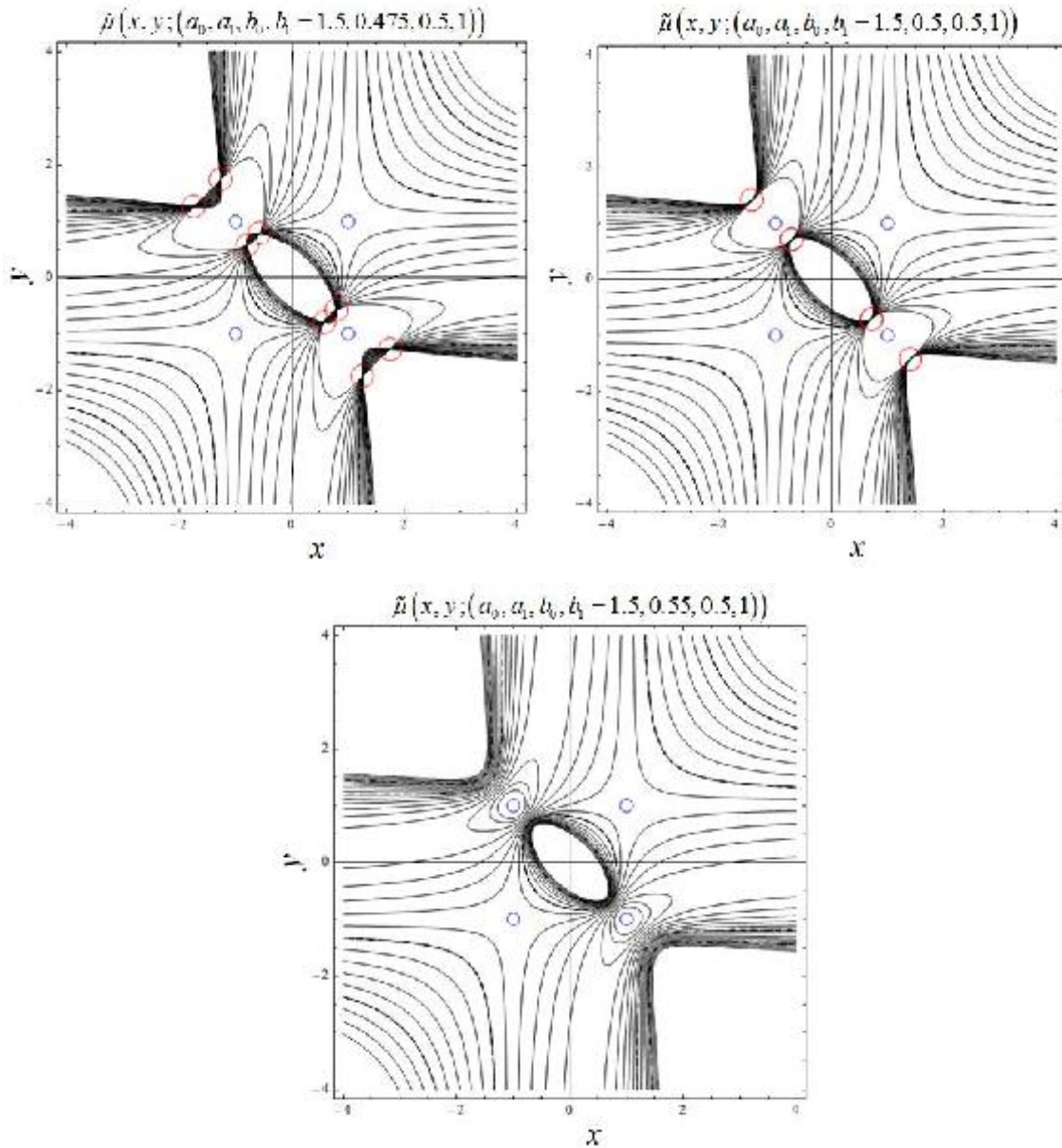


Figure 1. The level sets, the fixed points, the 2-periodic points, and the real base points of the integral (12) for parameter values  $a_0, a_1, b_0, b_1 = 1.5, 0.475, 0.5, 1$  (the diagram in the first row-left),  $a_0, a_1, b_0, b_1 = 1.5, 0.5, 0.5, 1$  (the diagram in the first row-right), and  $a_0, a_1, b_0, b_1 = 1.5, 0.55, 0.5, 1$  (the diagram in the second row).

$$y_{10} = \frac{x_{10}}{2} d \quad ;$$

$$x_{20} = \sqrt{\frac{\sqrt{d^2 - b_0 a_0 d^2 - 8 d - 8d^2 - 16 a_0 d - a_0 d^2 - b_0 - b_0 d}}{2 d d^2 2}} ;$$

$$y_{20} = \frac{x_{20}}{2} d \quad ; \quad d = \frac{b_1}{a_1}, a_1 = 0; \quad \sqrt{d^2 - 4}.$$

All base points depend on the parameters  $a_0, a_1, b_0$ , and  $b_1$ . To obtain the real base points for

$x_0, y_0 = x_{10}, y_{10}, \mp x_{10}, \mp y_{10}$ , the parameters should satisfy the following conditions

$$b_0 \in \mathbb{R} \quad a_0 = \frac{b_0}{d} \quad 2 \sqrt{\frac{2 d - 2d^2 - 4}{d^2 - d^2}} = 0 \quad d = 2 \quad 0 \quad d = \frac{2}{d} \quad d = 2 .$$

And then to obtain the real base points for  $x_0, y_0 = x_{20}, y_{20}, \mp x_{20}, \mp y_{20}$ , the parameters should satisfy the following conditions

$$b_0 \in \mathbb{R} \quad d = 2 \quad 0 \quad a_0 = \frac{b_0}{d} \quad 2 \sqrt{\frac{2 d - 2d^2 - 4}{d^2 - d^2}}$$

$$d = 2 \quad 0 \quad \frac{d^2 - 2}{d} \quad a_0 = \frac{b_0}{d} \quad 2 \sqrt{\frac{2 d - 2d^2 - 4}{d^2 - d^2}}$$

Figure 1 shows the base points (red-circle) together with fixed points and 2-periodic points (blue-circle). Three diagrams in Figure 1 are some level sets of integral (12) together with the fixed points, the 2-periodic points, and the base points for  $a_0 = 1.5$ . The diagram has fixed parameter values  $(a_1, b_0, b_1) = (0.475, 0.5, 1.0)$  (up-left diagram),  $(a_1, b_0, b_1) = (0.5, 0.5, 1.0)$  (up-right diagram) and  $(a_1, b_0, b_1) = (0.55, 0.5, 1.0)$  (down diagram).

### 4.3 Preimages of the Base Points

A base point is mapped to infinity by the mapping  $\hat{f}$ . Furthermore, there are two base points which do not refer to a base point. But after one iteration they are mapped to base points. This means that if we start at those points, the iteration of  $\hat{f}$  will be sent to infinity after two iterations. Let us name the set of points which are mapped by  $\hat{f}$  to a base point by  $P_1$ . We can then look at the set of points which are mapped by  $\hat{f}$  into  $P_1$ , or the preimage of  $P_1$  under  $\hat{f}$ ; and name the

set  $P_2$ . Continuing in a similar way, we constructed  $P_3, P_4$  and so on. Thus, the dynamical system (13) is well defined if we exclude the points in  $\bigcup_1 P_k$ .

Let us present a few explicit computations of  $P_k, k = 1, 2, 3$  as examples.

Consider the situation where  $(x, y) = \frac{1}{\sqrt{a_0}}, 0$  for  $a_0 > 0$  and denote the system (13) by

$$\hat{f} : \mathbb{R}^2 \rightarrow \mathbb{R}^2$$

$$(x, y) \mapsto \frac{a_1 y x^4 + 1 - x^3 b_0 - x + x^2 a_0 b_1}{a_1 x^4 + 1 - xy b_0 - xy x^2 - a_0 b_1}, x .$$

(16)

The preimage of  $(x, y) = \frac{1}{\sqrt{a_0}}, 0$  by (16) is the solution of the following system

$$x = 0$$

$$\frac{a_1 y x^4 + 1 - x^3 b_0 - x + x^2 a_0 b_1}{a_1 x^4 + 1 - xy b_0 - xy x^2 - a_0 b_1} = \frac{1}{\sqrt{a_0}} .$$

The solution is

$$P_1 = \hat{f}^{-1} \left( \frac{1}{\sqrt{a_0}}, 0 \right) = \left( 0, \frac{1}{\sqrt{a_0}} \right) \mid a_0 > 0 .$$

(17)

For the preimages  $P_2$  and  $P_3$ , we have

$$P_2 = \hat{f}^{-1} \left( 0, \frac{1}{\sqrt{a_0}} \right) = \left( \frac{3b_0}{4}, \frac{1}{\sqrt{a_0}} \right) \mid b_0 \in \mathbb{R} ,$$

$$P_3 = \hat{f}^{-1} \left( \frac{3b_0}{4}, \frac{1}{\sqrt{a_0}} \right) = \left( \frac{3b_0}{4}, \frac{1}{\sqrt{a_0}} \right) \mid \frac{4 - 1^2 - 4b_0^2}{4 - 1^2 - 8b_0^2} > 0, b_0 \in \mathbb{R} ,$$

where  $\frac{1}{\sqrt{a_0}} > 0$ .

The graphs of  $P_2$  to  $P_6$  for  $a_0 = 4$  are presented in Figure 2. The graph of  $P_2, P_3, P_4, P_5$ , and  $P_6$  are plotted using green, red, blue, orange, and black, respectively. It is interesting to note that if

we fix to one of the level sets of the integral, an preliminary observation shows an indication that there are only finitely many intersection points between  $\bigcup_1^6 P_k$ . Whether it is true when  $\bigcup_1 P_k$  is a subject of future investigation.

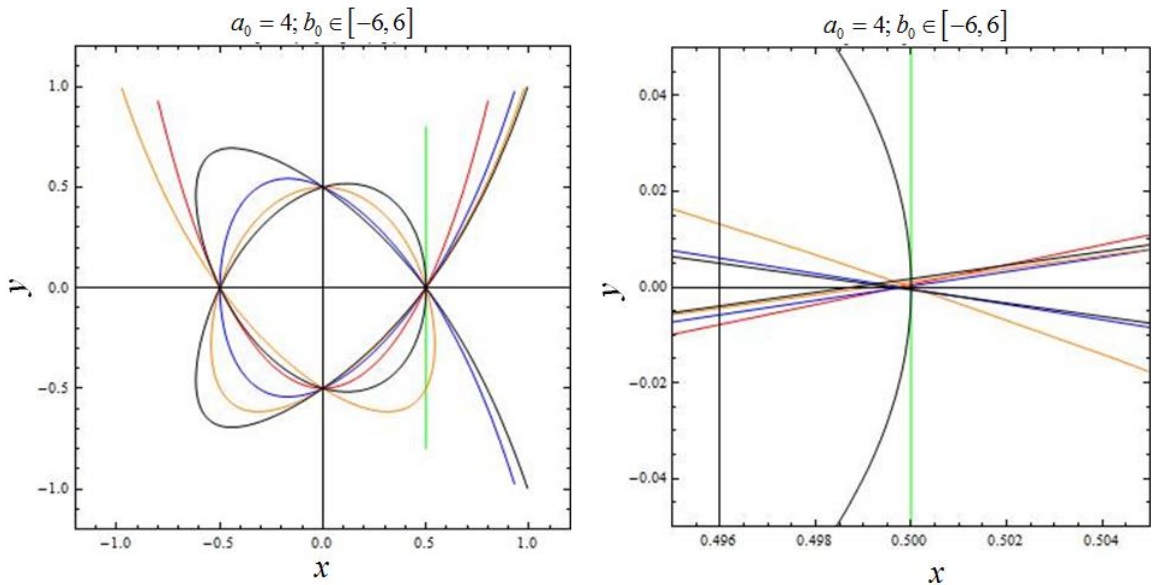


Figure 2. Plotting of the curves  $P_2 \dots P_6$  for  $a_0 = 4$ . We see that all curves pass through all base points (left). Enlarged view of the left image around a point  $(1/2, 0)$  to see the finer structure of the graph (right).

#### 4.4 Generic Situation

To study the dynamics of the system (13), we have plotted some of the level sets of the integral (see Figure 3). These level sets contain solution of the system (13). Furthermore, by studying how the level sets deformed as we vary the value parameter, the bifurcations in the system can be studied. This is however beyond the scope of this paper.

The two diagrams in the upper part of Figure 3 are some the level sets of integral (12) for  $\alpha_0 = 4$ . The diagram on the left-hand side has fixed  $(\alpha_1, \alpha_0, \alpha_1) = (0.7, 0.5, 1.0)$ , while for the diagram on the right-hand side is  $(\alpha_1, \alpha_0, \alpha_1) = (-0.875, 0.5, 1.0)$ . In the Figure 3 (down), we plot  $\alpha_0 = 4$ ,  $(\alpha_0, \alpha_1) = (0.5, 1.0)$ , and  $\alpha_1 = (-3, 0)$  around a fixed point  $(\alpha, \beta) = (-1, 1)$  (left); and we plot  $\alpha_0 = 4$ ,  $(\alpha_0, \alpha_1) = (0.5, 1.0)$ , and  $\alpha_1 = (-1, 2)$  around the fixed point  $(\alpha, \beta) = (1, 1)$  (right).

By setting  $\alpha_0 = 4$  and  $\alpha_1 = 0.7$ , we have all of the curves around the fixed points,  $(1, 1)$  and  $(-1, -1)$ , are of hyperbolic type. On the other hand we have all of the curves around the 2-periodic points,  $(-1, 1)$  and  $(1, -1)$ , are of elliptic type. But for  $\alpha_0 = 4$  and  $\alpha_1 = -0.875$ , the situation is reverse, i.e. all of the curves around the fixed points are of elliptic type while all of the curves around the 2-periodic points are of hyperbolic type.

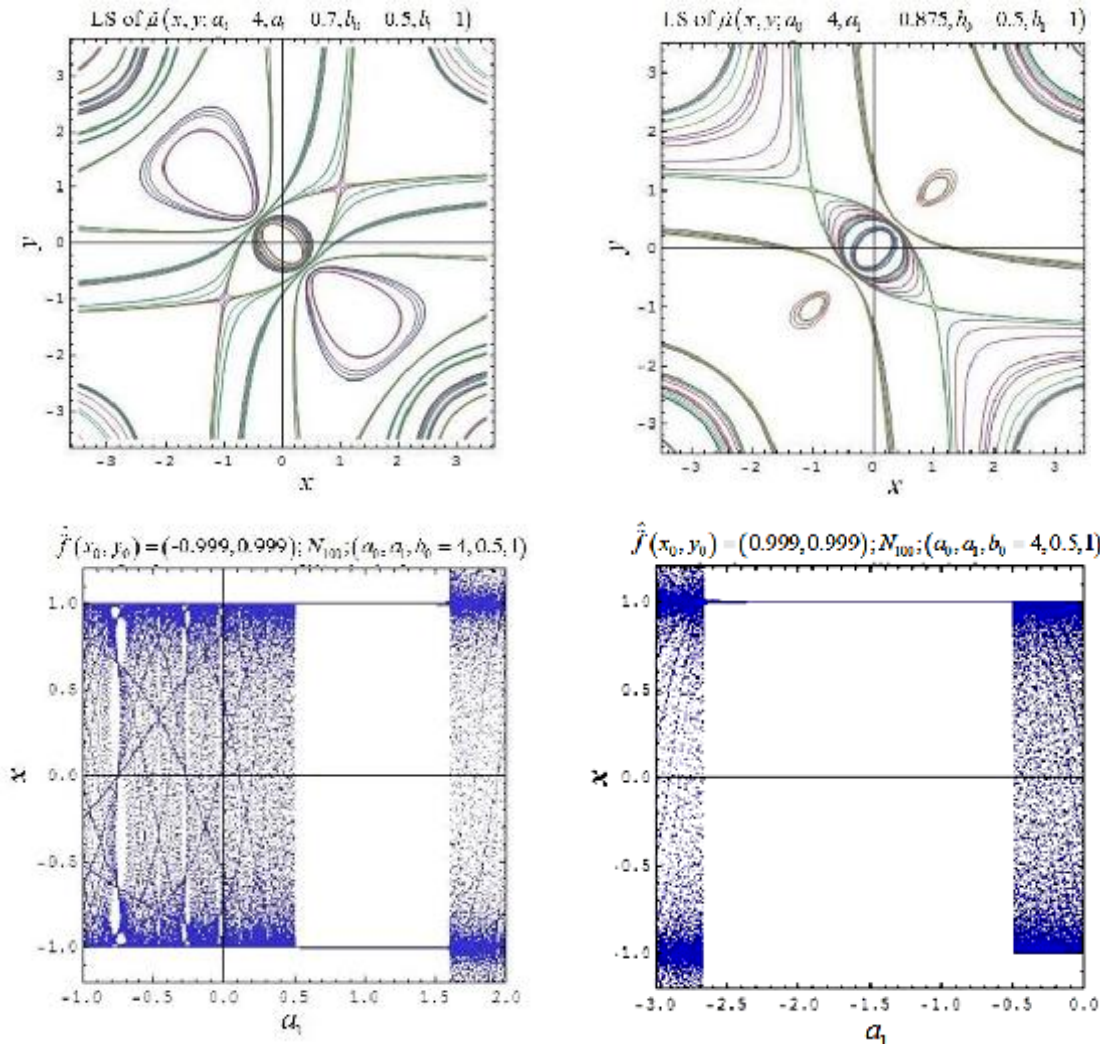


Figure 3. The level set forms of the integral (12) for the parameter value  $\alpha_0 = 4$  (up). Bifurcation situation corresponds with its (down).

#### 4.5 Stability

The stability of the system in the vicinity of the fixed points can be extracted from the integral function (see Kulenovic and Merino, 2002). In this case, we use the concept of Lyapunov stability to obtain the information of the stability of system (13). Note that the Hessian matrix evaluated at the fixed point  $(x^*, y^* = 1, 1)$  is

$$H = \begin{matrix} A & B \\ B & A \end{matrix}$$

where

$$A = \frac{2 a_1 b_0 a_0 + 1 b_1}{2 a_1 b_1^2};$$

$$B = \frac{2 a_1 + 4 b_0 + 1 a_0 b_1}{2 a_1 b_1^2}.$$

The determinant of the Hessian matrix  $H$  is

$$\text{Det } H = \frac{4 a_1 + 8 b_0 + 4 a_0 b_1}{2 a_1 b_1^3}$$

If the determinant of Hessian matrix  $H$  is positive then the integral (12) attains a minimum at  $x^*, y^* = 1, 1$ . In Table 1, we have listed the condition for the value of the parameters of the system, so that the determinant of  $H$  is positive. Consequently, the fixed points of (13),  $x^*, y^* = 1, 1$ , are the centre points (stable).

Table 1. The conditions for the value of the parameters so that the determinant of  $H$  is positive

No.	$a_0$	$b_1$	$b_0$	$a_1$
1.	$a_0 > 0$	$b_1 > 0$	$2 - 2a_0 < b_0 < 2$	$\frac{a_0 b_1}{2 b_0} < a_1 < \frac{b_1}{2}$
2.	$a_0 > 0$	$b_1 > 0$	$b_0 < 2 - 2a_0$	$\frac{b_1}{2} < a_1 < \frac{a_0 b_1}{2 b_0}$
3.	$a_0 > 0$	$b_1 > 0$	$b_0 > 2$	$a_1 < \frac{b_1}{2} \quad a_1 > \frac{a_0 b_1}{2 b_0}$
4.	$a_0 < 0$	$b_1 > 0$	$2 < b_0 - 2 - 2a_0$	$a_1 > \frac{b_1}{2} \quad a_1 < \frac{a_0 b_1}{2 b_0}$
5.	$a_0 < 0$	$b_1 > 0$	$b_0 > 2 - 2a_0$	$a_1 > \frac{a_0 b_1}{2 b_0} \quad a_1 < \frac{b_1}{2}$
6.	$a_0 < 0$	$b_1 > 0$	$b_0 < 2$	$\frac{b_1}{2} < a_1 < \frac{a_0 b_1}{2 b_0}$
7.	$a_0 > 0$	$b_1 < 0$	$2 - 2a_0 < b_0 < 2$	$\frac{b_1}{2} < a_1 < \frac{a_0 b_1}{2 b_0}$

Table 1 continued...

8.	$a_0 > 0$	$b_1 < 0$	$b_0 > 2$	$a_1 > \frac{b_1}{2}$
9.	$a_0 > 0$	$b_1 < 0$	$b_0 < 2 - 2a_0$	$\frac{a_0 b_1}{2 - b_0} < a_1 < \frac{b_1}{2}$
10.	$a_0 < 0$	$b_1 < 0$	$2 < b_0 - 2 - 2a_0$	$a_1 < \frac{b_1}{2} \quad a_1 > \frac{a_0 b_1}{2 - b_0}$
11.	$a_0 < 0$	$b_1 < 0$	$b_0 > 2 - 2a_0$	$a_1 > \frac{b_1}{2} \quad a_1 < \frac{a_0 b_1}{2 - b_0}$
12.	$a_0 < 0$	$b_1 < 0$	$b_0 < 2$	$\frac{a_0 b_1}{2 - b_0} < a_1 < \frac{b_1}{2}$
13.	$a_0 = 0$	$b_1 < 0$	$b_0 < 2$	$0 < a_1 < \frac{b_1}{2}$
14.	$a_0 = 0$	$b_1 < 0$	$b_0 > 2$	$a_1 > \frac{b_1}{2} \quad a_1 < 0$
15.	$a_0 = 0$	$b_1 > 0$	$b_0 < 2$	$\frac{b_1}{2} < a_1 < 0$
16.	$a_0 = 0$	$b_1 > 0$	$b_0 > 2$	$a_1 > 0 \quad a_1 < \frac{b_1}{2}$
17.	$a_0 \in \mathbb{R}$	$b_1 = 0$	$b_0 = 2$	$a_1 > 0$
18.	$a_0 < 0$	$b_1 = 0$	$b_0 \in \mathbb{R}$	$a_1 = 0$

## 5. Conclusion

After re-parametrizing the parameter in (9) and in (10) by  $a_0, a_1$  and  $\tilde{F} = \tilde{F}(b_0, b_1)$ , we have a new mapping  $\hat{f}$ . The properties of this mapping are integrable, measure preserving, and reversible. Furthermore, it has two fixed points and two 2-periodic points which are of an elliptic type and a hyperbolic type. It is interesting to underline the fact that we have the set of  $P_k$  consisting of points in  $\mathbb{R}^2$  which is mapped to infinity after  $k$ -iterates of the map.

Then we can consider the set of  $\mathcal{P} = \bigcup_1 P_k$ . For an arbitrary level set of the integral (12), the question should be whether the intersection between the set of  $\mathcal{P}$  with the level set is finite or infinite (could it be dense on the level set). This is a subject to future investigation.

### Conflict of Interest

The authors confirm that there is no conflict of interest for this publication.

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